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## LETTER TO THE EDITOR

# A Monte Carlo method for series expansions 

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#### Abstract

We describe a simple and general algorithm to calculate series expansions in enumeration problems to large orders approximately by a Monte Carlo method. It can be used to generate unbiased samples in cluster studies, e.g. linear or branched polymers, random surfaces, etc, in any dimension. We calculate the number of site animals of size $n$ on the square lattice for $n \leqslant 50$ and their average size to better than $1 \%$ accuracy.


The analysis of series expansions has been a very powerful technique in statistical mechanics. These calculations have been pushed to quite high orders for many problems, and in many cases yield the most accurate known estimates of physical quantities of interest, e.g. the critical exponents [1]. However, the computational effort needed to calculate the $n$ th-order term increases exponentially with $n$, and one is usually restricted to studying $n \leqslant 30$ in two dimensions [2]. In higher dimensions the situation is even worse, and for three-dimensional systems, series with more than 20 known terms are quite rare. Even if the asymptotic behaviour of the series is known theoretically in parametric form, precise estimates of the parameters such as critical exponents are hard to obtain from a short series because of significant corrections to scaling [3].

Often the calculation of series coefficients can be posed as an enumeration problem of some kinds of objects such as graphs or configurations (e.g. the linear or branched polymer problems). The exponential increase of the number of such configurations with their size $n$ makes a brute-force enumeration infeasible for large $n$.

In this letter we describe a simple algorithm, to be called the incomplete enumeration method, that allows an approximate determination of the number of configurations in enumeration problems by a Monte Carlo technique. The results obtained are only of finite accuracy, and are complementary to the exact series analysis because of the much higher orders reachable. While approximate values of high-order terms can be obtained by existing extrapolation methods such as exact series analysis or phenomenological renormalisation, it should be stressed that the error bars in these estimates are purely subjective. The magnitude of correction to the assumed extrapolation form is completely unknown. By contrast, the errors in the high-order coefficients calculated by the present method are purely statistical, objectively determined and can be made arbitrary small. Of course, extrapolated quantities such as critical parameters still have unknown, though smaller, systematic errors. The computation time increases only as a

[^0]power of $n$, and an accuracy of better than $1 \%$ is easily obtained for $n$ as large as 50 . The algorithm generates an unbiased sample [4], and other configuration averages can also be computed. It gives a direct estimate of the number of configurations, and hence of entropy.

As a test case, we studied the number of site animals $A_{n}$ on a square lattice, and their average moment of inertia $I_{n}$ as a function of the number of sites $n$ for $1 \leqslant n \leqslant 50$. An accuracy of about $0.5 \%$ for $A_{50}$ and about $0.13 \%$ for $I_{50}$ required approximately 20 h of cPU time on a CDC Cyber 175 machine. In contrast, the exact enumeration of $A_{24}$ already needs 10 months of CPU time on a PDP-11/70 computer [5] $\dagger$, and each successive order requires a four-fold increase in CPU time. On our machine, an exact calculation of $A_{50}$ would require approximately $4 \times 10^{20} \mathrm{~h}$.

For concreteness, we shall describe the algorithm for this special case. Other enumeration problems can be handled similarly. For the enumeration of lattice animals, there is a well known deterministic algorithm using backtracking [5,1], and an elegant and short fortran program using this algorithm has been published [6]. We describe it very briefly here.

One starts by choosing a rule for designating for each animal configuration one of its sites as the last added site. Deleting the last added site from an $n$-site animal $A$, we obtain an $(n-1)$-site animal $\boldsymbol{A}^{\prime}$, called the parent of $A$. The animal configurations are thus classified into a tree structure according to their lineage. The unique 1 -site animal forms the root of the tree. The $n$-site configurations are at a height ( $n-1$ ) connected to their parents at height $(n-2)$ by single bonds. The exhaustive enumeration of all animals having $n$ sites proceeds by a systematic exploration of the genealogical tree to height $(n-1)$ [6].

In the incomplete enumeration method, we arbitrarily choose a set of ( $n-1$ ) real numbers $p_{i}$ with $0<p_{i} \leqslant 1$ and $i=2$ to $n$. Any configuration with $r$ sites $(2 \leqslant r \leqslant n)$ is defined to be 'stillborn' with probability ( $1-p_{r}$ ) and normal with probability $p_{r}$, independent of the state of all other configurations. We delete the stillborn configurations from the genealogical tree. All configurations which are descendants of stillborn configurations are also deleted. We then systematically enumerate the remaining $n$-site configurations in the genealogical tree. This can be done very efficiently using backtracking. When a particular $r$-site configuration is first generated, one chooses to ignore it and all its descendants in the enumeration with probability ( $1-p_{r}$ ).

The probability that a particular $r$-site animal will be enumerated in a given trial is $p_{2} p_{3} \ldots p_{r} \stackrel{\text { def }}{=} \Pi_{r}$ (say), and is the same for all configurations with the same $r$. The algorithm thus generates an unbiased sample of configurations. The number of $n$-site animals generated in one trial is a random variable $X_{n}$. Averaging over several trials, we can estimate $\left\langle X_{n}\right\rangle$. Since this equals $\pi_{n} A_{n}$, we can estimate $A_{n}$ using Monte Carlo methods.

By judiciously choosing $p_{i}$, we can get $\left\langle X_{n}\right\rangle$ to be close to 1 . Then the percolation process on the genealogical tree is close to threshold, and the variance of $X_{n}$ increases slowly with $n$, and may be approximated as increasing like $n^{a}$, where $a$ is some exponent. For a genealogical tree with constant branching number, $a$ can be calculated analytically, and one finds $a=1$. For trees with variable branchings, $a$ is harder to calculate exactly. To obtain the fractional error in $A_{n}$ less than $\varepsilon$, the number of trials has to increase as $\varepsilon^{-2} n^{a}$ for large $n$. Since the time required to generate an animal

[^1]increases linearly with $n$, the computer time to obtain a fractional error in $A_{n}$ increases as $\varepsilon^{-2} n^{a+1}$.

The method may be viewed as constructing a site percolation process on the genealogical tree, with the stillborn configurations as the blocked nodes on the tree. The enumerated configurations form the connected part of the tree.

It is useful to contrast the Monte Carlo algorithm outlined here with those commonly used. The conventional method, as developed by Metropolis et al [7] relies on the construction of a Markov process whose time averages replace the ensemble averages sought. Averaging over the initial state may also be done [8]. This procedure gives rise to strong time correlations in samples and, in many problems of physical interest, the approach to equilibrium distribution can become very slow due to hydrodynamic slowing, or critical slowing or metastability [8]. Also it is not possible to determine thermodynamic quantities such as entropy or free energy directly from the simulations and these must be calculated by integration over temperatures.

The present algorithm does not construct a Markov process, and the problem of correlations is much more manageable. Different configurations produced in a single trial are, of course, correlated with each other being more likely to share a common lineage. There is also a small anticorrelation since a configuration can occur at most once in a single trial (sampling without replacement).

The variability of $X_{n}$ in different trials is crucial in keeping the sample unbiased, as is the case in the simple algorithm for generating self-avoiding walks by randomly generating an ideal random walk of $n$ steps, and rejecting it if it is found to be self-intersecting. This algorithm is inefficient for large $n$, as the fraction of accepted configurations is very small. In the incomplete enumeration approach discussed here, a step that would lead to rejection makes the algorithm backtrack to attempt another allowed completion, and does not lead to a rejection of the full configuration.

The basic idea of combining an enumeration algorithm with indeterminism has been discussed earlier by Redner and Reynolds [9] in the case of self-avoiding walks and by Yang et al [10]. For the animals problem a similar algorithm was proposed by Lam [11]. However, the detailed algorithm used by Lam is not completely free of bias. The present algorithm differs from his by having a different rule for backtracking. His algorithm may try to generate some configurations more than once during the execution of the program (see figure 1). This makes the sample biased in favour of such configurations. Another source of bias in it is introduced by terminating a trial before completion as soon as some preset number of configurations has been generated.

The approach is easily generalised to other enumeration problems such as linear or branched polymers in restricted geometries, random surfaces, etc. In each case the exponential growth of the number of configurations is curtailed by randomly pruning off branches in the genealogical tree of configurations. The technique can also be


Figure 1. An example to illustrate the difference between the present algorithm and that given in [11]. We show two animals of size 4 on a square lattice. By Lam's algorithm [11], the second animal may be generated by sequential occupation of sites abcdor acdb or $a b d c$, and thus is three times as likely to be generated in a single trial as the first animal. This bias is avoided in the present algorithm.
generalised to generate configurations with prescribed (unequal) weights (needed, for example, in study systems like the Ising model in the canonical ensemble). It is not suitable for determining the behaviour of non-leading singularities. A detailed analysis of the algorithm and additional numerical results will be published elsewhere.

For the numerical study of site animals on the square lattice, we chose $p_{r}=(r+1) / 4 r$ for $r \geqslant 2$, and averaged over $4 \times 10^{7}$ independent trials. The mean number $\left\langle X_{n}\right\rangle$ obtained for $n=3-50$ is shown in table 1 . The standard error of estimate shown was obtained by grouping the data into 150 equal parts and calculating fluctuations about the mean value. The average moment of inertia $I_{n}$ was calculated in a separate run of $4 \times 10^{6}$ trials.

Table 1. The mean number of animals generated per trial $\left\langle X_{n}\right\rangle$ as a function of $n$. The estimated total number of $n$-site animals is $A_{n}=\left\langle X_{n}\right\rangle 4^{n-1} 2 /(n+1)$.

| $n$ | $\left\langle X_{n}\right\rangle$ | $n$ | $\left\langle X_{n}\right\rangle$ | $n$ | $\left\langle X_{n}\right\rangle$ |
| ---: | :--- | :--- | :--- | :--- | :--- |
| 3 | $0.7497 \pm 0.0004$ | 19 | $0.8641 \pm 0.0019$ | 35 | $1.0943 \pm 0.0040$ |
| 4 | $0.7423 \pm 0.0005$ | 20 | $0.8765 \pm 0.0020$ | 36 | $1.1120 \pm 0.0043$ |
| 5 | $0.7383 \pm 0.0006$ | 21 | $0.8888 \pm 0.0022$ | 37 | $1.1287 \pm 0.0044$ |
| 6 | $0.7382 \pm 0.0007$ | 22 | $0.9019 \pm 0.0023$ | 38 | $1.1461 \pm 0.0046$ |
| 7 | $0.7423 \pm 0.0008$ | 23 | $0.9158 \pm 0.0023$ | 39 | $1.1632 \pm 0.0048$ |
| 8 | $0.7479 \pm 0.0008$ | 24 | $0.9292 \pm 0.0024$ | 40 | $1.1812 \pm 0.0050$ |
| 9 | $0.7559 \pm 0.0009$ | 25 | $0.9427 \pm 0.0026$ | 41 | $1.1997 \pm 0.0052$ |
| 10 | $0.7645 \pm 0.0010$ | 26 | $0.9568 \pm 0.0027$ | 42 | $1.2185 \pm 0.0054$ |
| 11 | $0.7741 \pm 0.0011$ | 27 | $0.9704 \pm 0.0029$ | 43 | $1.2368 \pm 0.0056$ |
| 12 | $0.7844 \pm 0.0012$ | 28 | $0.9853 \pm 0.0030$ | 44 | $1.2561 \pm 0.0058$ |
| 13 | $0.7944 \pm 0.0013$ | 29 | $0.9997 \pm 0.0031$ | 45 | $1.2753 \pm 0.0060$ |
| 14 | $0.8053 \pm 0.0014$ | 30 | $1.0149 \pm 0.0032$ | 46 | $1.2950 \pm 0.0062$ |
| 15 | $0.8160 \pm 0.0015$ | 31 | $1.0295 \pm 0.0034$ | 47 | $1.3146 \pm 0.0065$ |
| 16 | $0.8275 \pm 0.0016$ | 32 | $1.0454 \pm 0.0035$ | 48 | $1.3344 \pm 0.0068$ |
| 17 | $0.8393 \pm 0.0018$ | 33 | $1.0611 \pm 0.0037$ | 49 | $1.3549 \pm 0.0071$ |
| 18 | $0.8515 \pm 0.0019$ | 34 | $1.0775 \pm 0.0038$ | 50 | $1.3759 \pm 0.0074$ |

The values $\left\langle X_{n}\right\rangle$ vary approximately exponentially with $n$ for $n \geqslant 15$, and a linear plot of $\log \left\langle X_{n}\right\rangle$ against $n$ gives

$$
\ln \left\langle X_{n}\right\rangle \simeq \alpha n+\beta
$$

with $\alpha=0.0153 \pm 0.0004$ and $\beta=-0.455 \pm 0.020$. No term proportional to $\log n$ is expected, as the exactly known exponent $\theta=1$ has been incorporated in our choice of $p_{n}$. These values agree with the result $\alpha=0.01552 \pm 0.00004$ obtained by phenomenological renormalisation [12], and $\alpha=0.01553 \pm 0.00005, \beta=-0.456 \pm 0.010$ obtained by extrapolation of the exact series coefficients [13]†. A plot of $n \log \left(I_{n+1} / I_{n}\right)$ against $1 / n$ is an approximate straight line which extrapolated to $n \rightarrow \infty$ gives $I_{n} \sim n^{2 \nu+1}$ with $\nu=0.646 \pm 0.004$. This value of $\nu$ is also in agreement with earlier, more precise estimates. The advantage of the present technique is that it can be used with equal ease in higher dimensions, where the existing estimates are much less precise [14].

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[^2]
## References

[1] Domb C and Green M S (ed) 1974 Phase Transitions and Critical Phenomena vol 3 (New York: Academic)
[2] Baxter R J, Enting I G and Tsang S K 1980 J. Stat. Phys. 22465
[3] Adler J, Moshe M and Privman V 1983 Ann. Israel Phys. Soc. 5397
[4] Glaus U 1985 J. Phys. A: Math. Gen. 18 L609
Beretti A and Sokal A D 1985 J. Stat. Phys. 40483
Aragao de Carvalho C and Carcibolo S 1983 J. Physique 44323
[5] Redelmeier D H 1981 Discrete Math. 36191
Sykes M F 1986 J. Phys. A: Math. Gen. 191007
[6] Redner S 1982 J. Stat. Phys. 29309
Demme E S and Diemer K 1984 J. Undergrad. Res. in Phys. 325
[7] Metropolis N, Rosenbluth A W, Rosenbluth M N, Teller A H and Teller E 1953 J. Chem. Phys. 211087
[8] Binder K (ed) 1979 Monte Carlo Methods (Berlin: Springer)
[9] Redner S and Reynolds P J 1981 J. Phys. A: Math. Gen. 142679 Fisher M E, Privman V and Redner S 1984 J. Phys. A: Math. Gen. 17 L569
Privman V and Redner S 1985 J. Phys. A: Math. Gen. 18 L781
[10] Yang Y S, Liu Y and Lam P M 1985 Z. Phys. B 59445
[11] Lam P M 1986 J. Phys. A: Math. Gen. 19 L155
[12] Derrida B and Stauffer D 1985 J. Physique 461623 Kertész J 1986 J. Phys. A: Math. Gen. 19599
[13] Guttman A J 1982 J. Phys. A: Math. Gen. 151987
[14] Stanley H E, Redner S and Yang Z R 1982 J. Phys. A: Math. Gen. 15 L569


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[^1]:    $\dagger$ The computer time could be decreased considerably by using a suitable generating function technique (see Sykes [5]).

[^2]:    $\dagger$ Our parameter $\alpha=-\log \left(4 x_{\mathrm{c}}\right)$ is in the notation of the authors in [12].

